## Nabla in Curvilinear Coordinates

Reference: M. R. Spiegel, Schaum's Outline of ... Vector Analysis ..., Chapter 7 (and part of Chap. 8). (Page references are to that book.)

Suppose that $\vec{A}=A_{r} \hat{r}+A_{\theta} \hat{\theta}+A_{\phi} \hat{\phi}$ with respect to the usual basis of unit vectors in spherical coordinates. What is $\nabla \cdot \vec{A}$ ? It is not $\frac{\partial A_{r}}{\partial r}+$ $\frac{\partial A_{\theta}}{\partial \theta}+\frac{\partial A_{\phi}}{\partial \phi}$. The reason is that the spherical unit vectors themselves are functions of position, and their own derivatives must be taken into account. The correct formula is [p. 165, solution 63(b)]

$$
\nabla \cdot \vec{A}=\frac{1}{r^{2}} \frac{\partial}{\partial r}\left(r^{2} A_{r}\right)+\frac{1}{r \sin \theta} \frac{\partial}{\partial \theta}\left(\sin \theta A_{\theta}\right)+\frac{1}{r \sin \theta} \frac{\partial A_{\phi}}{\partial \phi} .
$$

Where did this come from?
To set the stage, let's recall that the most elementary local basis associated with a coordinate system is the tangent vectors

$$
\left\{\frac{\partial \vec{r}}{\partial u_{j}}\right\} \quad\left[\text { in the spherical case, } \quad\left\{\frac{\partial \vec{r}}{\partial r}, \frac{\partial \vec{r}}{\partial \theta}, \frac{\partial \vec{r}}{\partial \phi}\right\}\right] .
$$

The metric tensor is defined as the matrix of inner products

$$
g_{i j}=\frac{\partial \vec{r}}{\partial u_{i}} \cdot \frac{\partial \vec{r}}{\partial u_{j}}
$$

This means that the formula for arc length is

$$
d s^{2}=\sum_{i j} g_{i j} d u_{i} d u_{j} \quad\left[=d r^{3}+r^{2} d \theta^{2}+r^{2} \sin ^{2} \theta d \phi^{2} \quad \text { in sphericals }\right] .
$$

(In this context $d v d u=+d u d v$, not $-d u d v$ as in surface integrals.)

Orthogonal coordinates are the very important special cases where $g$ is a diagonal matrix:

$$
g=\left(\begin{array}{ccc}
h_{1}{ }^{2} & 0 & 0 \\
0 & h_{2}{ }^{2} & 0 \\
0 & 0 & h_{3}{ }^{2}
\end{array}\right) \quad \text { (in three-dimensional cases). }
$$

That is, by definition

$$
h_{j}=\left\|\frac{\partial \vec{r}}{\partial u_{j}}\right\| ;
$$

in the spherical case,

$$
h_{r}=1, \quad h_{\theta}=r, \quad h_{\phi}=r \sin \theta .
$$

In orthogonal coordinates we can define the orthonormal basis of unit vectors,

$$
\hat{u}_{j}=\frac{1}{h_{j}} \frac{\partial \vec{r}}{\partial u_{j}} .
$$

The other standard basis vectors, the normal vectors $\nabla u_{j}$, satisfy

$$
\hat{u}_{j}=h_{j} \nabla u_{j} .
$$

Now we turn to the gradient. Note that there is some ambiguity in the term "components of the gradient in spherical coordinates", because we have three natural spherical bases to choose among. The most useful choice for "everyday" physics is the orthonormal basis, which is obtained by a rotation from the Cartesian basis $\{\hat{\imath}, \hat{\jmath}, \hat{k}\}$. By the chain rule one sees, for instance, that

$$
\frac{\partial f}{\partial \theta}=\sum_{j} \frac{\partial f}{\partial x_{j}} \frac{\partial x_{j}}{\partial \theta}=\nabla f \cdot \frac{\partial \vec{r}}{\partial \theta}=h_{\theta} \hat{\theta} \cdot \nabla f .
$$

But since this basis is ON, this is equivalent to saying that the $\hat{\theta}$ component of $\nabla f$ is $h_{\theta}{ }^{-1} \frac{\partial f}{\partial \theta}$. In general [p. 137],

$$
\nabla f=\sum_{j} \frac{1}{h_{j}} \frac{\partial f}{\partial u_{j}} \hat{u}_{j},
$$

which in the spherical case is [p. 165, solution 63(a)]

$$
\nabla f=\frac{\partial f}{\partial r} \hat{r}+\frac{1}{r} \frac{\partial f}{\partial \theta} \hat{\theta}+\frac{1}{r \sin \theta} \frac{\partial f}{\partial \phi} \hat{\phi}
$$

To treat the divergence and curl, first note that

$$
\hat{u}_{1}=\hat{u}_{2} \times \hat{u}_{3}=h_{2} h_{3}\left(\nabla u_{2}\right) \times\left(\nabla u_{3}\right)
$$

(and the two obvious cyclic permutations of this formula). Now consider one term of the divergence:

$$
\begin{aligned}
\nabla \cdot\left(A_{1} \hat{u}_{1}\right) & =\nabla \cdot\left(A_{1} h_{2} h_{3} \nabla u_{2} \times \nabla u_{3}\right) \\
& =\nabla\left(A_{1} h_{2} h_{3}\right) \cdot\left(\nabla u_{2} \times \nabla u_{3}\right)+A_{1} h_{2} h_{3} \nabla \cdot\left(\nabla u_{2} \times \nabla u_{3}\right) .
\end{aligned}
$$

But the last term is $\overrightarrow{0}$, because $\nabla \cdot(\vec{B} \times \vec{C})=\vec{C} \cdot(\nabla \times \vec{B})-\vec{B}(\nabla \times \vec{C})$ and $\nabla \times \nabla u_{j}=0$. So

$$
\begin{aligned}
\nabla \cdot\left(A_{1} \hat{u}_{1}\right) & =\nabla\left(A_{1} h_{2} h_{3}\right) \cdot\left(\nabla u_{2} \times \nabla u_{3}\right) \\
& =\frac{1}{h_{2} h_{3}}\left(\hat{u}_{2} \times \hat{u}_{3}\right) \cdot \nabla\left(A_{1} h_{1} h_{2}\right) \\
& =\frac{\hat{u}_{1}}{h_{1} h_{3}} \cdot\left[\frac{\partial}{\partial u_{1}}\left(A_{1} h_{2} h_{3}\right) \cdot \frac{1}{h_{1}} \hat{u}_{1}+\left(\hat{u}_{2} \text { and } \hat{u}_{3} \text { terms }\right)\right],
\end{aligned}
$$

where the previous formula for the gradient has been used. Only the $\hat{u}_{1}$ term survives the dot product, so

$$
\nabla \cdot\left(A_{1} \hat{u}_{1}\right)=\frac{1}{h_{1} h_{2} h_{3}} \frac{\partial}{\partial u_{1}}\left(A_{1} h_{2} h_{3}\right) .
$$

Adding the two cyclic analogues of this term one gets [p. 150]

$$
\nabla \cdot \vec{A}=\frac{1}{h_{1} h_{2} h_{3}}\left[\frac{\partial}{\partial u_{1}}\left(A_{1} h_{2} h_{3}\right)+\frac{\partial}{\partial u_{2}}\left(A_{2} h_{3} h_{1}\right)+\frac{\partial}{\partial u_{3}}\left(A_{3} h_{1} h_{2}\right)\right] .
$$

In spherical coordinates this reduces to the formula given at the beginning.
Incidentally, the denominator in the divergence formula, $h_{1} h_{2} h_{3}=$ $\sqrt{\operatorname{det} g}$, is equal to the Jacobian determinant that occurs in the curvilinear
volume formula. This fact is a consequence of Part 2 of the Volume Theorem on p. 340 of Linearity, where the metric tensor was introduced without a name.

Similarly [p. 150; p. 154 for spherical case], the curl is

$$
\nabla \times \vec{A}=\frac{1}{h_{1} h_{2} h_{3}}\left\|\begin{array}{ccc}
h_{1} \hat{u}_{1} & h_{2} \hat{u}_{2} & h_{3} \hat{u}_{3} \\
\frac{\partial}{\partial u_{1}} & \frac{\partial}{\partial u_{2}} & \frac{\partial}{\partial u_{3}} \\
h_{1} A_{1} & h_{2} A_{2} & h_{3} A_{3}
\end{array}\right\| .
$$

An alternative way of deriving the divergence and curl formulas [pp. $151-152$ ] is based on the integral definitions of divergence and curl [p. 365 of Linearity]. The factors $h_{j}$ show up in the relationships between coordinate increments $\left(\Delta u_{j}\right)$ and physical lengths, areas, and volumes.

