## Nabla in Curvilinear Coordinates

**Reference:** M. R. Spiegel, *Schaum's Outline of* ... *Vector Analysis* ... , Chapter 7 (and part of Chap. 8). (Page references are to that book.)

Suppose that  $\vec{A} = A_r \hat{r} + A_\theta \hat{\theta} + A_\phi \hat{\phi}$  with respect to the usual basis of unit vectors in spherical coordinates. What is  $\nabla \cdot \vec{A}$ ? It is not  $\frac{\partial A_r}{\partial r} + \frac{\partial A_\theta}{\partial \theta} + \frac{\partial A_\phi}{\partial \phi}$ . The reason is that the spherical unit vectors themselves are functions of position, and their own derivatives must be taken into account. The correct formula is [p. 165, solution 63(b)]

$$\nabla \cdot \vec{A} = \frac{1}{r^2} \frac{\partial}{\partial r} (r^2 A_r) + \frac{1}{r \sin \theta} \frac{\partial}{\partial \theta} (\sin \theta A_\theta) + \frac{1}{r \sin \theta} \frac{\partial A_\phi}{\partial \phi}$$

Where did this come from?

To set the stage, let's recall that the most elementary local basis associated with a coordinate system is the tangent vectors

$$\left\{\frac{\partial \vec{r}}{\partial u_j}\right\} \qquad \left[\text{in the spherical case,} \quad \left\{\frac{\partial \vec{r}}{\partial r}, \frac{\partial \vec{r}}{\partial \theta}, \frac{\partial \vec{r}}{\partial \phi}\right\}\right].$$

The metric tensor is defined as the matrix of inner products

$$g_{ij} = \frac{\partial \vec{r}}{\partial u_i} \cdot \frac{\partial \vec{r}}{\partial u_j}$$

This means that the formula for arc length is

$$ds^{2} = \sum_{ij} g_{ij} du_{i} du_{j} \qquad \left[ = dr^{3} + r^{2} d\theta^{2} + r^{2} \sin^{2} \theta d\phi^{2} \quad \text{in sphericals} \right].$$

(In this context dv du = +du dv, not -du dv as in surface integrals.)

Orthogonal coordinates are the very important special cases where g is a diagonal matrix:

$$g = \begin{pmatrix} h_1^2 & 0 & 0\\ 0 & h_2^2 & 0\\ 0 & 0 & h_3^2 \end{pmatrix}$$
 (in three-dimensional cases).

That is, by definition

$$h_j = \left\| \frac{\partial \vec{r}}{\partial u_j} \right\|;$$

in the spherical case,

$$h_r = 1, \quad h_\theta = r, \quad h_\phi = r \sin \theta.$$

In orthogonal coordinates we can define the orthonormal basis of unit vectors,

$$\hat{u}_j = \frac{1}{h_j} \frac{\partial \vec{r}}{\partial u_j} \,.$$

The other standard basis vectors, the normal vectors  $\nabla u_j$ , satisfy

$$\hat{u}_j = h_j \nabla u_j$$
.

Now we turn to the gradient. Note that there is some ambiguity in the term "components of the gradient in spherical coordinates", because we have three natural spherical bases to choose among. The most useful choice for "everyday" physics is the orthonormal basis, which is obtained by a rotation from the Cartesian basis  $\{\hat{i}, \hat{j}, \hat{k}\}$ . By the chain rule one sees, for instance, that

$$\frac{\partial f}{\partial \theta} = \sum_{j} \frac{\partial f}{\partial x_{j}} \frac{\partial x_{j}}{\partial \theta} = \nabla f \cdot \frac{\partial \vec{r}}{\partial \theta} = h_{\theta} \hat{\theta} \cdot \nabla f.$$

But since this basis is ON, this is equivalent to saying that the  $\hat{\theta}$  component of  $\nabla f$  is  $h_{\theta}^{-1} \frac{\partial f}{\partial \theta}$ . In general [p. 137],

$$\nabla f = \sum_{j} \frac{1}{h_j} \frac{\partial f}{\partial u_j} \hat{u}_j \,,$$

which in the spherical case is [p. 165, solution 63(a)]

$$\nabla f = \frac{\partial f}{\partial r}\,\hat{r} + \frac{1}{r}\,\frac{\partial f}{\partial \theta}\,\hat{\theta} + \frac{1}{r\sin\theta}\,\frac{\partial f}{\partial \phi}\,\hat{\phi}.$$

To treat the *divergence* and curl, first note that

$$\hat{u}_1 = \hat{u}_2 \times \hat{u}_3 = h_2 h_3 (\nabla u_2) \times (\nabla u_3)$$

(and the two obvious cyclic permutations of this formula). Now consider one term of the divergence:

$$\nabla \cdot (A_1 \hat{u}_1) = \nabla \cdot (A_1 h_2 h_3 \nabla u_2 \times \nabla u_3)$$
  
=  $\nabla (A_1 h_2 h_3) \cdot (\nabla u_2 \times \nabla u_3) + A_1 h_2 h_3 \nabla \cdot (\nabla u_2 \times \nabla u_3).$ 

But the last term is  $\vec{0}$ , because  $\nabla \cdot (\vec{B} \times \vec{C}) = \vec{C} \cdot (\nabla \times \vec{B}) - \vec{B}(\nabla \times \vec{C})$  and  $\nabla \times \nabla u_j = 0$ . So

$$\nabla \cdot (A_1 \hat{u}_1) = \nabla (A_1 h_2 h_3) \cdot (\nabla u_2 \times \nabla u_3)$$
  
=  $\frac{1}{h_2 h_3} (\hat{u}_2 \times \hat{u}_3) \cdot \nabla (A_1 h_1 h_2)$   
=  $\frac{\hat{u}_1}{h_1 h_3} \cdot \left[ \frac{\partial}{\partial u_1} (A_1 h_2 h_3) \cdot \frac{1}{h_1} \hat{u}_1 + (\hat{u}_2 \text{ and } \hat{u}_3 \text{ terms}) \right],$ 

where the previous formula for the gradient has been used. Only the  $\hat{u}_1$  term survives the dot product, so

$$\nabla \cdot (A_1 \hat{u}_1) = \frac{1}{h_1 h_2 h_3} \frac{\partial}{\partial u_1} (A_1 h_2 h_3).$$

Adding the two cyclic analogues of this term one gets [p. 150]

$$\nabla \cdot \vec{A} = \frac{1}{h_1 h_2 h_3} \left[ \frac{\partial}{\partial u_1} (A_1 h_2 h_3) + \frac{\partial}{\partial u_2} (A_2 h_3 h_1) + \frac{\partial}{\partial u_3} (A_3 h_1 h_2) \right].$$

In spherical coordinates this reduces to the formula given at the beginning.

Incidentally, the denominator in the divergence formula,  $h_1h_2h_3 = \sqrt{\det g}$ , is equal to the Jacobian determinant that occurs in the curvilinear

volume formula. This fact is a consequence of Part 2 of the Volume Theorem on p. 340 of *Linearity*, where the metric tensor was introduced without a name.

Similarly [p. 150; p. 154 for spherical case], the curl is

$$\nabla \times \vec{A} = \frac{1}{h_1 h_2 h_3} \left\| \begin{array}{ccc} h_1 \hat{u}_1 & h_2 \hat{u}_2 & h_3 \hat{u}_3 \\ \frac{\partial}{\partial u_1} & \frac{\partial}{\partial u_2} & \frac{\partial}{\partial u_3} \\ h_1 A_1 & h_2 A_2 & h_3 A_3 \end{array} \right\|.$$

An alternative way of deriving the divergence and curl formulas [pp. 151–152] is based on the integral definitions of divergence and curl [p. 365 of *Linearity*]. The factors  $h_j$  show up in the relationships between coordinate increments  $(\Delta u_j)$  and physical lengths, areas, and volumes.